

# Optical Constants of Water in the Infrared

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The results of our earlier studies of reflection and absorption in various spectral regions are reviewed and then used to provide values of the complex index of refraction  $\hat{N} = n + ik$  of water at 27°C in the spectral range 5000–10 cm<sup>-1</sup>, corresponding to wavelengths in the range 2 μm to 1 mm. Values of  $n$ ,  $k$ , and the Lambert absorption coefficient  $\alpha$ , which are presented graphically and in tabular form, should prove useful in studies of the scattering of infrared radiation by water droplets in the atmosphere and in studies of radiative heat balance at water surfaces.

Although the infrared spectrum of water had been the subject of numerous investigations, *Irvine and Pollack* [1968] made a critical survey of published results that revealed many inconsistencies and a general paucity of quantitative data on which to base values of the real and imaginary parts of the complex index of refraction  $\hat{N} = n + ik$ . In view of the importance of  $n$  and  $k$  in calculations of the transmission, scattering, and absorption of electromagnetic radiation by water droplets in the earth's atmosphere, our laboratory group has devoted considerable attention to the quantitative determination of the optical properties of water in the infrared. We have based our earlier listings of the optical constants  $n$  and  $k$  on quantitative measurements of various types in various spectral regions. The purpose of the present paper is to give a critical review of our earlier studies with the purpose of providing a set of 'best values' for use in atmospheric studies.

In our initial study, covering the 5000- to 400-cm<sup>-1</sup> region, *Querry et al.* [1969] attempted to measure the reflectance of polarized radiation at two large angles of incidence and to determine  $n$  and  $k$  by solution of the generalized Fresnel equations. In the range 5000–330 cm<sup>-1</sup>, *Rusk et al.* [1971] employed reflectance measurements at near-normal incidence and at an angle of 53° near Brewster's angle. Although there was fair agreement in the values of  $n$  and  $k$  obtained in these two studies, serious uncertainties were introduced as a result of the imperfect polarizers employed at nonnormal incidence in the first study and the failure to achieve Brewster's angle in the vicinity of absorption bands in the second study.

In view of these uncertainties, *Hale et al.* [1972] applied a Kramers-Kronig (KK) phase shift analysis to obtain values of the optical constants from *Rusk's* measurements of reflectance at near-normal incidence. The values of  $n$  based on the KK analysis represented an improvement on the earlier values; the KK analysis gave good values of  $k$  in the vicinity of strong absorption maximums but was unreliable in spectral regions where  $k$  is small. In general, reflectance measurements can give reliable values for  $n$  and also for large  $k$ ; they thus complement careful absorption measurements, which can provide reliable values for small  $k$  but somewhat questionable values for large  $k$ .

Our next study by *Robertson and Williams* [1971] was the quantitative measurement of the Lambert absorption coefficient  $\alpha$  defined by  $I = I_0 \exp[-\alpha x]$ ; in this work we used a wedge-shaped absorption cell designed by *Robertson*, and we covered the spectral range 4300–300 cm<sup>-1</sup>. The values of  $k = \lambda\alpha/4\pi$  based on absorption measurements were more precise than those based on reflectance in spectral regions of

small  $k$  and agreed, within the stated limits of uncertainty, in the centers of absorption bands where  $k$  is large; in the spectral range  $\nu < 600$  cm<sup>-1</sup> the uncertainties in  $k$  became larger because of limitations imposed by the spectrometers employed. The values of  $k$  were measured in this low-frequency region by *Robertson et al.* [1973], who used a far infrared grating instrument to determine  $\alpha$  in the spectral range between 800 and 50 cm<sup>-1</sup>; these authors also obtained values of  $n$  by means of a KK analysis of measured values of  $\alpha$ .

In the spectroscopy of the remote infrared, interferometers used with Fourier transform techniques have marked advantages over conventional grating instruments. Using interferometric methods, *John Chamberlain* and his associates at the National Physical Laboratory (NPL) have obtained values of  $n$  and  $k$  in the range 100–20 cm<sup>-1</sup>; in the course of this work, *Davies et al.* [1970] employed absorption techniques, and *Zafar et al.* [1973] employed reflection techniques. Existing water data in the microwave and radiofrequency regions have been summarized by *Ray* [1972].

## PRESENT STUDY

In preparing the present summary of our work on water we have based our values of the optical constants primarily on (1) *Robertson's* absorption measurements and (2) *Rusk's* measurements of spectral reflectance at near-normal incidence. In extensions of these primary data to the near-infrared and visible we have made use of the recent work of *Palmer and Williams* [1974]; in the extreme infrared we have used the NPL results in the 100- to 20-cm<sup>-1</sup> region and results taken from *Ray's* survey in the frequency range below 20 cm<sup>-1</sup>. In spectral regions where accurate values of absorption coefficients and reflectances have been determined independently we have obtained values of  $n$  and  $k$  from Fresnel's equation; in other regions we have employed KK methods.

The refractive index  $n$  can be determined from the KK relation

$$n(\nu) = 1 + (1/2\pi^2)P \int_0^\infty \frac{\alpha(\nu') - \alpha(\nu)}{\nu'^2 - \nu^2} d\nu' \quad (1)$$

where  $\alpha$  represents the Lambert absorption coefficient, for which we have values in the range between the radiofrequency region and 14,500 cm<sup>-1</sup> in the visible. In order to obtain values of  $n$  in the infrared from (1) it is sufficient to take account of ultraviolet contributions by assuming a single far ultraviolet band which will give the proper value of  $n$  at some frequency for which it is accurately known from independent measurements; we chose characteristics for the hypothetical ultraviolet band that would yield a value  $n = 1.306$  at 5000 cm<sup>-1</sup> in agreement with all our own earlier measurements.

On the basis of  $n$  evaluated from (1) and of direct experimental values of  $k = \lambda\alpha/4\pi$  we calculated the values of the normal incidence reflectance  $R$  in the range 800–120  $\text{cm}^{-1}$ ; these calculated values of  $R$  served to check Rusk's values in the 800- to 350- $\text{cm}^{-1}$  range and to provide values of  $R$  in the 350- to 120- $\text{cm}^{-1}$  range, where no reflectance measurements have been made. In the 120- to 90- $\text{cm}^{-1}$  range we joined our calculated values to a reflectance curve for the 90- to 10- $\text{cm}^{-1}$  range calculated from the NPL optical constants and those listed by Ray. On the basis of measured and calculated values of reflectance over the whole range from the near ultraviolet to the radiofrequency range, we then employed the KK phase shift theorem

$$\phi(\nu) = (2\nu/\pi)P \int_0^\infty \frac{\ln [R(\nu')^{1/2}]}{\nu'^2 - \nu^2} d\nu' \quad (2)$$

where  $[R(\nu)]^{1/2}$  is the modulus of the complex reflectivity  $\hat{R} = [R(\nu)]^{1/2} \exp [i\phi(\nu)]$ . In terms of  $\phi$  and  $R$  the values of  $n$  and  $k$  at any frequency are given by the relations

$$n = (1 - R)/(1 + R - 2R^{1/2} \cos \phi) \quad (3)$$

$$k = (-2R^{1/2} \sin \phi)/(1 + R - 2R^{1/2} \cos \phi) \quad (4)$$

We have used (2) along with (3) and (4) to provide  $n$  and  $k$  over the entire frequency range of present interest.

In the computer programs used for the solution of (1) and (2) we have employed methods based on Simpson's rule with a basic increment of 10  $\text{cm}^{-1}$  except in the vicinity of the singularity at  $\nu$ , where analytic solutions involving quadratic approximations of  $\alpha(\nu)$  and  $\ln [R(\nu)]^{1/2}$  were used. The 10- $\text{cm}^{-1}$  mesh is satisfactory over most of the range of present interest but becomes coarse at the lowest frequencies.

OPTICAL CONSTANTS

In Figure 1 we give our final values of the absorption index  $k$  as a function of frequency in waves per centimeter and wavelength in micrometers. The values represent the weighted average of  $k$  based on direct measurements of  $\alpha$  and on KK analyses; greater weight is given to the values based on direct measurement. The error bars shown in the figure represent the maximum differences between measured values and values based on (1), (2), and (4); the error bars thus give a measure of the in-

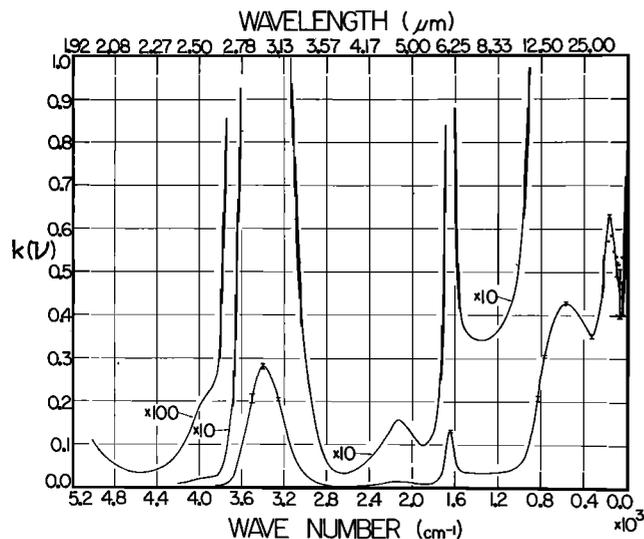


Fig. 1. Absorption index  $k$  as a function of wave number and wavelength.

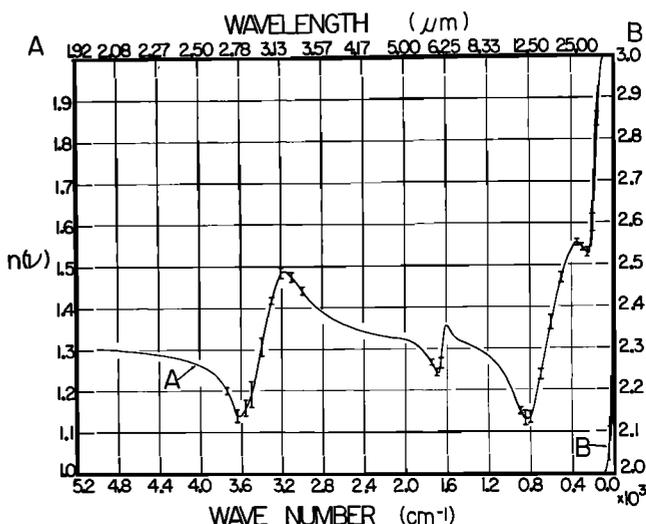


Fig. 2. Refractive index  $n$  as a function of wave number and wavelength.

ternal consistency of our work. In general, the actual uncertainties, which have been estimated in our earlier papers, are comparable with those given by the error bars except in the vicinity of the strong absorption band near 3400  $\text{cm}^{-1}$ , for which Robertson and Williams [1971] list an uncertainty of  $\pm 4\%$  in  $k$ ; thus since transmission measurements tend to give an underestimate of  $k$  at the centers of strong absorption bands, our results indicate that  $k$  may be as large as 0.294 at 3390  $\text{cm}^{-1}$ .

Our final values of the refractive index  $n$  are plotted as a function of wave number and wavelength in Figure 2. The curve shown represents a weighted average of direct determinations in regions where  $\alpha$  and  $R$  have been determined directly by experiment, of KK determinations from (1), and of KK determinations from (2) and (3); greater weight has been accorded to direct determinations. It should be noted that our values of  $n$  in the range 350–120  $\text{cm}^{-1}$  are based entirely on

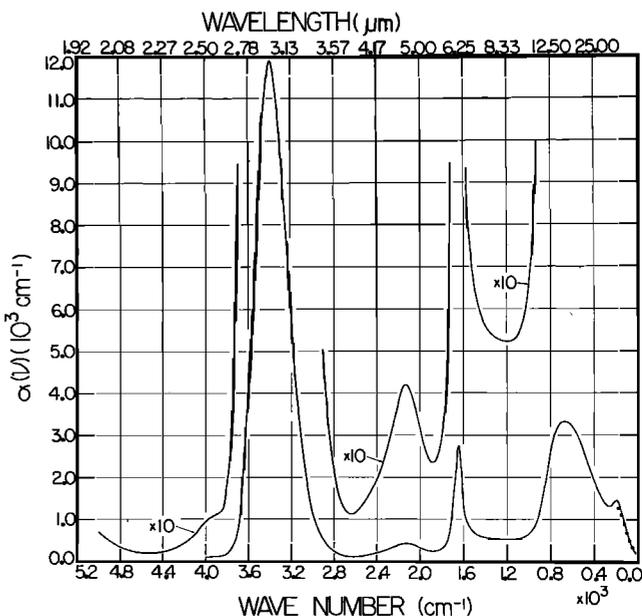


Fig. 3. Lambert absorption coefficient  $\alpha$  as a function of wave number and wavelength. Values shown are based on direct measurement of absorption.

KK analyses, since we have made no reflection measurements in this range. The error bars on the curve in the figure are a measure of the self-consistency of our work and at a given frequency represent the maximum differences between  $n$  as determined by various techniques; the large uncertainties in spectral regions where  $n$  is changing rapidly are probably due in part to spectrometer calibration problems and in part to the size of the increments employed in the KK analysis.

TABLE 1. (continued)

$\nu$	$n(\nu)$	$k(\nu)$	$\alpha(\nu)$	$\lambda$	$\nu$	$n(\nu)$	$k(\nu)$	$\alpha(\nu)$	$\lambda$
5000	1.303	0.00110	69.2	2.000	3500	1.191	0.206	8750.0	2.857
4950	1.301	0.000900	56.0	2.020	3490	1.199	0.218	9270.0	2.865
4900	1.301	0.000731	45.0	2.041	3480	1.212	0.229	9660.0	2.874
4850	1.300	0.000617	37.6	2.062	3470	1.220	0.239	10120.0	2.882
4800	1.298	0.000514	31.0	2.083	3460	1.233	0.249	10500.0	2.890
4750	1.298	0.000452	27.0	2.105	3450	1.246	0.258	10850.0	2.899
4700	1.296	0.000400	23.6	2.128	3440	1.258	0.265	11150.0	2.907
4650	1.295	0.000359	21.0	2.151	3430	1.271	0.271	11370.0	2.915
4600	1.294	0.000341	19.7	2.174	3420	1.282	0.276	11600.0	2.924
4550	1.293	0.000338	19.3	2.198	3410	1.293	0.280	11780.0	2.933
4500	1.291	0.000345	19.5	2.222	3400	1.305	0.281	11850.0	2.941
4450	1.289	0.000376	21.0	2.247	3390	1.317	0.282	11900.0	2.950
4400	1.287	0.000416	23.0	2.273	3380	1.329	0.282	11870.0	2.959
4350	1.285	0.000465	25.4	2.299	3370	1.342	0.279	11720.0	2.967
4300	1.282	0.000542	29.3	2.326	3360	1.353	0.276	11600.0	2.976
4250	1.280	0.000652	34.8	2.353	3350	1.364	0.272	11400.0	2.985
4200	1.277	0.000792	41.8	2.381	3340	1.376	0.267	11150.0	2.994
4150	1.274	0.000968	50.5	2.410	3330	1.386	0.262	10920.0	3.003
4100	1.270	0.00123	63.5	2.439	3320	1.398	0.255	10570.0	3.012
4050	1.265	0.00156	79.5	2.469	3310	1.407	0.250	10300.0	3.021
4000	1.261	0.00190	95.7	2.500	3300	1.417	0.243	10000.0	3.030
3990	1.260	0.00195	97.5	2.506	3290	1.426	0.236	9670.0	3.040
3980	1.259	0.00200	100.0	2.513	3280	1.434	0.228	9300.0	3.049
3970	1.257	0.00205	102.0	2.519	3270	1.442	0.220	8950.0	3.058
3960	1.256	0.00207	103.0	2.525	3260	1.450	0.212	8570.0	3.067
3950	1.255	0.00210	104.0	2.532	3250	1.457	0.204	8270.0	3.077
3940	1.254	0.00212	105.0	2.538	3240	1.465	0.195	7820.0	3.086
3930	1.252	0.00215	106.0	2.545	3230	1.471	0.183	7320.0	3.096
3920	1.250	0.00219	108.0	2.551	3220	1.476	0.173	6830.0	3.106
3910	1.249	0.00224	110.0	2.558	3210	1.480	0.163	6400.0	3.115
3900	1.247	0.00227	111.0	2.564	3170	1.487	0.125	4840.0	3.155
3890	1.246	0.00231	113.0	2.571	3160	1.487	0.117	4550.0	3.165
3880	1.243	0.00234	114.0	2.577	3150	1.486	0.110	4320.0	3.175
3870	1.241	0.00239	116.0	2.584	3140	1.485	0.0994	3890.0	3.185
3860	1.240	0.00243	118.0	2.591	3130	1.482	0.0920	3620.0	3.195
3850	1.238	0.00248	120.0	2.597	3120	1.479	0.0855	3390.0	3.205
3840	1.235	0.00257	124.0	2.604	3110	1.477	0.0785	3120.0	3.215
3830	1.232	0.00270	130.0	2.611	3100	1.474	0.0716	2840.0	3.226
3820	1.230	0.00298	143.0	2.618	3090	1.472	0.0653	2590.0	3.236
3810	1.227	0.00330	158.0	2.625	3080	1.467	0.0600	2390.0	3.247
3800	1.224	0.00402	192.0	2.632	3070	1.464	0.0550	2190.0	3.257
3790	1.221	0.00437	208.0	2.639	3060	1.461	0.0504	2010.0	3.268
3780	1.218	0.00482	229.0	2.646	3050	1.457	0.0462	1840.0	3.279
3770	1.214	0.00536	254.0	2.653	3040	1.454	0.0422	1680.0	3.289
3760	1.210	0.00627	296.0	2.660	3030	1.451	0.0385	1530.0	3.300
3750	1.205	0.00732	345.0	2.667	3020	1.448	0.0348	1390.0	3.311
3740	1.200	0.00855	402.0	2.674	3010	1.444	0.0315	1260.0	3.322
3730	1.195	0.0105	490.0	2.681	3000	1.441	0.0297	1120.0	3.333
3720	1.191	0.0127	593.0	2.688	2990	1.437	0.0279	1050.0	3.344
3710	1.185	0.0145	677.0	2.695	2980	1.434	0.0262	980.0	3.356
3700	1.179	0.0164	762.0	2.703	2970	1.431	0.0250	933.0	3.367
3690	1.172	0.0186	862.0	2.710	2960	1.427	0.0229	850.0	3.378
3680	1.166	0.0205	946.0	2.717	2950	1.425	0.0210	780.0	3.390
3670	1.157	0.0282	1300.0	2.725	2940	1.421	0.0193	713.0	3.401
3660	1.149	0.0380	1930.0	2.732	2930	1.418	0.0177	650.0	3.413
3650	1.144	0.0462	2270.0	2.740	2920	1.415	0.0163	599.0	3.425
3640	1.139	0.0548	2600.0	2.747	2910	1.413	0.0151	551.0	3.436
3630	1.138	0.0649	2970.0	2.755	2900	1.410	0.0138	503.0	3.448
3620	1.138	0.0744	3340.0	2.762	2890	1.407	0.0128	466.0	3.460
3610	1.139	0.0836	3720.0	2.770					

TABLE 1. (continued)

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$\nu$	$n(\nu)$	$k(\nu)$	$\alpha(\nu)$	$\lambda$	$\nu$	$n(\nu)$	$k(\nu)$	$\alpha(\nu)$	$\lambda$
2880	1.405	0.0118	428.0	3.472	2180	1.327	0.0145	396.0	4.587
2870	1.403	0.0110	398.0	3.484	2170	1.327	0.0149	406.0	4.608
2860	1.400	0.0101	363.0	3.497	2160	1.327	0.0152	412.0	4.630
2850	1.398	0.00941	337.0	3.509	2150	1.327	0.0154	417.0	4.651
2840	1.396	0.00866	309.0	3.521	2140	1.326	0.0156	419.0	4.673
2830	1.394	0.00807	287.0	3.534	2130	1.326	0.0157	419.0	4.695
2820	1.392	0.00737	261.0	3.546	2120	1.326	0.0157	418.0	4.717
2810	1.390	0.00683	241.0	3.559	2110	1.325	0.0157	416.0	4.739
2800	1.388	0.00625	220.0	3.571	2100	1.325	0.0155	410.0	4.762
2790	1.387	0.00579	203.0	3.584	2090	1.325	0.0153	402.0	4.785
2780	1.385	0.00538	188.0	3.597	2080	1.325	0.0151	394.0	4.808
2770	1.383	0.00506	176.0	3.610	2070	1.325	0.0148	386.0	4.831
2760	1.382	0.00473	164.0	3.623	2060	1.325	0.0146	377.0	4.854
2750	1.379	0.00449	155.0	3.636	2050	1.324	0.0143	368.0	4.878
2740	1.378	0.00424	146.0	3.650	2040	1.324	0.0140	359.0	4.902
2730	1.377	0.00405	139.0	3.663	2030	1.323	0.0137	349.0	4.926
2720	1.375	0.00389	133.0	3.676	2020	1.322	0.0133	338.0	4.950
2710	1.374	0.00376	128.0	3.690	2010	1.322	0.0129	327.0	4.975
2700	1.372	0.00363	123.0	3.704	2000	1.321	0.0126	317.0	5.000
2690	1.371	0.00355	120.0	3.717	1990	1.320	0.0122	306.0	5.025
2680	1.370	0.00347	117.0	3.731	1980	1.319	0.0118	294.0	5.051
2670	1.369	0.00340	114.0	3.745	1970	1.318	0.0115	284.0	5.076
2660	1.367	0.00335	112.0	3.759	1960	1.318	0.0110	272.0	5.102
2650	1.366	0.00336	112.0	3.774	1950	1.317	0.0108	264.0	5.128
2640	1.365	0.00335	111.0	3.788	1940	1.316	0.0105	255.0	5.155
2630	1.363	0.00339	112.0	3.802	1930	1.314	0.0103	249.0	5.181
2620	1.361	0.00340	112.0	3.817	1920	1.313	0.0101	244.0	5.208
2610	1.361	0.00348	114.0	3.831	1910	1.311	0.0100	240.0	5.236
2600	1.360	0.00352	115.0	3.846	1900	1.310	0.00993	237.0	5.263
2590	1.358	0.00363	118.0	3.861	1890	1.308	0.00990	235.0	5.291
2580	1.358	0.00370	120.0	3.876	1880	1.306	0.00995	235.0	5.319
2570	1.357	0.00378	122.0	3.891	1870	1.304	0.0100	236.0	5.348
2560	1.355	0.00389	125.0	3.906	1860	1.302	0.0102	238.0	5.376
2550	1.354	0.00399	128.0	3.922	1850	1.299	0.0104	242.0	5.405
2540	1.353	0.00410	131.0	3.937	1840	1.297	0.0107	247.0	5.435
2530	1.352	0.00422	134.0	3.953	1830	1.294	0.0110	253.0	5.464
2520	1.351	0.00433	137.0	3.968	1820	1.291	0.0115	262.0	5.495
2510	1.350	0.00450	142.0	3.984	1810	1.288	0.0120	274.0	5.525
2500	1.349	0.00465	146.0	4.000	1800	1.285	0.0128	289.0	5.556
2490	1.348	0.00479	150.0	4.016	1790	1.282	0.0138	311.0	5.587
2480	1.348	0.00494	154.0	4.032	1780	1.278	0.0150	336.0	5.618
2470	1.347	0.00512	159.0	4.049	1770	1.275	0.0166	370.0	5.650
2460	1.346	0.00531	164.0	4.065	1760	1.271	0.0185	409.0	5.682
2450	1.345	0.00549	169.0	4.082	1750	1.267	0.0205	451.0	5.714
2440	1.344	0.00568	174.0	4.098	1740	1.262	0.0242	529.0	5.747
2430	1.344	0.00586	179.0	4.115	1730	1.256	0.0293	637.0	5.780
2420	1.343	0.00608	185.0	4.132	1720	1.251	0.0332	734.0	5.814
2410	1.342	0.00631	191.0	4.149	1710	1.247	0.0429	947.0	5.848
2400	1.341	0.00653	197.0	4.167	1700	1.242	0.0544	1200.0	5.882
2390	1.340	0.00673	202.0	4.184	1690	1.241	0.0688	1515.0	5.917
2380	1.340	0.00696	208.0	4.202	1680	1.241	0.0840	1840.0	5.952
2370	1.338	0.00722	215.0	4.219	1670	1.247	0.1021	2175.0	5.988
2360	1.337	0.00749	222.0	4.237	1660	1.265	0.117	2430.0	6.024
2350	1.337	0.00779	230.0	4.255	1650	1.289	0.130	2670.0	6.061
2340	1.335	0.00806	237.0	4.274	1640	1.311	0.132	2738.0	6.098
2330	1.334	0.00833	244.0	4.292	1630	1.332	0.124	2566.0	6.135
2320	1.334	0.00864	252.0	4.310	1620	1.349	0.106	2139.0	6.173
2310	1.333	0.00896	260.0	4.329	1610	1.354	0.0880	1760.0	6.211
2300	1.332	0.00927	268.0	4.348	1600	1.356	0.0740	1465.0	6.250
2290	1.332	0.00966	278.0	4.367	1590	1.354	0.0618	1200.0	6.289
2280	1.331	0.0100	287.0	4.386	1580	1.350	0.0535	1025.0	6.329
2270	1.330	0.0104	297.0	4.405	1570	1.345	0.0484	934.0	6.369
2260	1.330	0.0108	308.0	4.425	1560	1.341	0.0447	863.0	6.410
2250	1.330	0.0112	318.0	4.444	1550	1.337	0.0420	806.0	6.452
2240	1.329	0.0117	330.0	4.464	1540	1.333	0.0398	758.0	6.494
2230	1.329	0.0122	342.0	4.484	1530	1.330	0.0383	726.0	6.536
2220	1.329	0.0126	352.0	4.505	1520	1.326	0.0373	703.0	6.579
2210	1.328	0.0131	364.0	4.525	1510	1.324	0.0370	683.0	6.623
2200	1.328	0.0136	376.0	4.545	1500	1.322	0.0366	666.0	6.667
2190	1.327	0.0140	386.0	4.566					

TABLE 1. (continued)

TABLE 1. (continued)

$\nu$	$n(\nu)$	$k(\nu)$	$\alpha(\nu)$	$\lambda$	$\nu$	$n(\nu)$	$k(\nu)$	$\alpha(\nu)$	$\lambda$
1490	1.320	0.0363	652.0	6.711	780	1.142	0.292	2883.0	12.821
1480	1.319	0.0360	638.0	6.757	770	1.157	0.305	2969.0	12.987
1470	1.318	0.0357	624.0	6.803	760	1.171	0.317	3040.0	13.158
1460	1.317	0.0355	612.0	6.849	750	1.182	0.328	3100.0	13.333
1450	1.316	0.0352	602.0	6.897	740	1.189	0.338	3150.0	13.514
1440	1.315	0.0350	593.0	6.944	730	1.201	0.347	3192.0	13.699
1430	1.314	0.0347	582.0	6.993	720	1.213	0.356	3231.0	13.889
1420	1.313	0.0346	575.0	7.042	710	1.223	0.365	3263.0	14.085
1410	1.311	0.0343	564.0	7.092					
1400	1.310	0.0342	558.0	7.143	700	1.236	0.373	3287.0	14.286
1390	1.309	0.0342	554.0	7.194	690	1.249	0.379	3298.0	14.493
1380	1.308	0.0342	550.0	7.246	680	1.264	0.386	3307.0	14.706
1370	1.307	0.0343	547.0	7.299	670	1.277	0.392	3308.0	14.925
1360	1.306	0.0342	543.0	7.353	660	1.289	0.397	3307.0	15.152
1350	1.305	0.0342	540.0	7.407	650	1.303	0.403	3301.0	15.385
1340	1.303	0.0342	537.0	7.463	640	1.313	0.408	3291.0	15.625
1330	1.302	0.0342	535.0	7.519	630	1.324	0.412	3276.0	15.873
1320	1.301	0.0342	532.0	7.576	620	1.335	0.417	3259.0	16.129
1310	1.300	0.0344	530.0	7.634	610	1.348	0.420	3234.0	16.393
1300	1.298	0.0345	530.0	7.692	600	1.361	0.423	3203.0	16.667
1290	1.296	0.0346	529.0	7.752	590	1.372	0.425	3167.0	16.949
1280	1.295	0.0349	528.0	7.813	580	1.385	0.427	3126.0	17.241
1270	1.294	0.0351	527.0	7.874	570	1.396	0.428	3077.0	17.544
1260	1.293	0.0351	526.0	7.937	560	1.407	0.427	3022.0	17.857
1250	1.291	0.0351	525.0	8.000	550	1.419	0.427	2964.0	18.182
1240	1.288	0.0352	524.0	8.065	540	1.431	0.426	2903.0	18.519
1230	1.286	0.0356	524.0	8.130	530	1.441	0.425	2842.0	18.868
1220	1.285	0.0359	523.0	8.197	520	1.451	0.423	2779.0	19.231
1210	1.283	0.0361	523.0	8.264	510	1.462	0.421	2709.0	19.608
1200	1.281	0.0362	522.0	8.333	500	1.470	0.418	2638.0	20.000
1190	1.279	0.0366	522.0	8.403	490	1.480	0.415	2565.0	20.408
1180	1.276	0.0370	523.0	8.475	480	1.488	0.411	2494.0	20.833
1170	1.274	0.0374	523.0	8.547	470	1.496	0.408	2423.0	21.277
1160	1.271	0.0378	523.0	8.621	460	1.504	0.404	2347.0	21.739
1150	1.269	0.0383	524.0	8.696	450	1.510	0.401	2280.0	22.222
1140	1.267	0.0387	525.0	8.772	440	1.515	0.397	2210.0	22.727
1130	1.264	0.0392	527.0	8.850	430	1.521	0.394	2143.0	23.256
1120	1.261	0.0398	529.0	8.929	420	1.527	0.390	2072.0	23.810
1110	1.259	0.0405	532.0	9.009	410	1.532	0.386	2004.0	24.390
1100	1.256	0.0411	536.0	9.091	400	1.537	0.382	1933.0	25.000
1090	1.253	0.0417	540.0	9.174	390	1.541	0.377	1862.0	25.641
1080	1.249	0.0424	546.0	9.259	380	1.545	0.372	1793.0	26.316
1070	1.246	0.0434	553.0	9.346	370	1.549	0.368	1724.0	27.027
1060	1.242	0.0443	561.0	9.434	360	1.552	0.363	1658.0	27.778
1050	1.238	0.0453	571.0	9.524	350	1.552	0.359	1593.0	28.571
1040	1.234	0.0467	583.0	9.615	340	1.552	0.356	1532.0	29.412
1030	1.230	0.0481	596.0	9.709	330	1.550	0.352	1472.0	30.303
1020	1.224	0.0497	613.0	9.804	320	1.546	0.353	1432.0	31.250
1010	1.220	0.515	631.0	9.901	310	1.543	0.357	1401.0	32.258
1000	1.214	0.0534	651.0	10.000	300	1.541	0.361	1374.0	33.333
990	1.208	0.0557	673.0	10.101	290	1.539	0.368	1351.0	34.483
980	1.202	0.0589	702.0	10.204	280	1.537	0.375	1331.0	35.714
970	1.194	0.0622	733.0	10.309	270	1.534	0.385	1317.0	37.037
960	1.189	0.0661	770.0	10.417	260	1.532	0.398	1311.0	38.462
950	1.181	0.0707	817.0	10.526	250	1.529	0.414	1310.0	40.000
940	1.174	0.0764	866.0	10.638	240	1.525	0.436	1323.0	41.667
930	1.168	0.0828	927.0	10.753	230	1.528	0.469	1364.0	43.478
920	1.162	0.0898	993.0	10.870	220	1.542	0.505	1407.0	45.455
910	1.156	0.0973	1064.0	10.989	210	1.567	0.539	1434.0	47.619
900	1.149	0.107	1165.0	11.111	200	1.600	0.571	1445.0	50.000
890	1.143	0.118	1270.0	11.236	190	1.640	0.597	1437.0	52.632
880	1.139	0.130	1396.0	11.364	180	1.689	0.618	1412.0	55.556
870	1.135	0.144	1533.0	11.494	170	1.746	0.629	1358.0	58.824
860	1.132	0.159	1682.0	11.628	160	1.801	0.622	1266.0	62.500
850	1.132	0.176	1833.0	11.765	150	1.848	0.608	1165.0	66.667
840	1.131	0.192	1987.0	11.905	140	1.890	0.593	1065.0	71.429
830	1.132	0.208	2143.0	12.048	130	1.929	0.577	967.0	76.923
820	1.130	0.226	2309.0	12.195	120	1.960	0.557	872.0	83.333
810	1.130	0.243	2467.0	12.346	110	1.982	0.532	773.0	90.909
800	1.134	0.260	2618.0	12.500	100	1.997	0.507	678.0	100.000
790	1.138	0.277	2760.0	12.658	90	2.000	0.487	594.0	111.111
					80	2.010	0.466	509.0	125.000

TABLE 1. (continued)

$\nu$	$n(\nu)$	$k(\nu)$	$\alpha(\nu)$	$\lambda$
70	2.020	0.450	429.0	142.857
60	2.040	0.444	360.0	166.667
50	2.070	0.438	290.0	200.000
40	2.110	0.460	240.0	250.000
30	2.150	0.527	210.0	333.333
20	2.225	0.718	192.0	500.000
10	2.600	1.0902	137.0	1000.000

Frequencies  $\nu$  are expressed in waves per centimeter ( $\text{cm}^{-1}$ ),  $n(\nu)$  and imaginary  $k(\nu)$  parts of the dielectric constant  $N = n + ik$  are dimensionless, the Lambert absorption coefficient  $\alpha(\nu)$  defined by the relation  $I = I_0 \exp[-\alpha x]$  is expressed in waves per centimeter ( $\text{cm}^{-1}$ ), and wavelengths  $\lambda$  are given in micrometers ( $\mu\text{m}$ ). The values of  $n(\nu)$ ,  $k(\nu)$ , and  $\alpha(\nu)$  are given for water at 27°C.

Although the values of  $n$  and  $k$  provide all the information actually required for a quantitative description of the optical properties of water, a set of values of the Lambert coefficient  $\alpha$  is of direct use in providing information of importance to studies of radiative heat balance at horizontal water surfaces. We have therefore included a plot of  $\alpha$  versus  $\nu$  in Figure 3; values of  $\alpha$  given in the plot are based entirely on direct measurements and thus differ slightly from  $\alpha$  values calculated from our averaged values of  $k$ . The values of  $\alpha$  given in Figure 3 would apply in good approximation to clear freshwater lakes and can provide rough approximations of the properties of clear seawater, as suggested by *Hobson and Williams* [1971].

Irvine and Pollack have emphasized the importance of presenting optical constants in tabular as well as graphical form. In Table 1 we list our best values of  $k$ ,  $n$ , and  $\alpha$  at frequency intervals of  $10 \text{ cm}^{-1}$  over most of the range between 5000 and  $10 \text{ cm}^{-1}$ ; the values in the table correspond to those plotted in Figures 1-3 and involve the same uncertainties. These values apply to water at a laboratory temperature of approximately 27°C; values at other temperatures can be estimated from the plots given by *Hale et al.* [1972]. A molecular interpretation of the water spectrum was given by *Robertson et al.* [1973].

#### COMPARISON WITH OTHER STUDIES

Our values for the optical constants can be compared with those obtained in earlier studies by *Pontier and Dechambenoy* [1965, 1966] in France and by *Zolotarev et al.* [1969] in Russia. The present results for  $k$  are in excellent agreement with both of these studies in the range 5000-4000  $\text{cm}^{-1}$  but are in somewhat serious disagreement in the vicinity of the strong absorption band near 3400  $\text{cm}^{-1}$ , where plots of the earlier studies differ by several percent from those in Figure 1. The peak values of  $k$  in the two studies are 0.305, a value somewhat higher than our present highest estimate and 8% higher than the value given in our plot; the absorption band obtained by the French workers is centered at a slightly lower frequency than the frequency given in the other studies.

In the frequency range 2800-800  $\text{cm}^{-1}$  there is truly excellent agreement between the present  $k$  values and those reported by the Russian group; throughout most of this region the French values of  $k$  are significantly higher than our values. At frequencies lower than 800  $\text{cm}^{-1}$  the French values are generally

greater than ours, and the Russian values generally lower; through the entire region below 800  $\text{cm}^{-1}$  the  $k$  values reported by the other groups fall within  $\pm 10\%$  of the values we give in Figure 1.

In comparing our present values for  $n$  with the earlier studies we find that in the 5000- to 3600- $\text{cm}^{-1}$  region our values are in good agreement with the values obtained in the French study; throughout this region the Russian values are considerably lower than ours and are in serious disagreement in the 3800- to 3600- $\text{cm}^{-1}$  range, where the Russian values are much lower than ours. In the range 3200-400  $\text{cm}^{-1}$  the  $n$  values obtained in the earlier studies generally fall within  $\pm 1\%$  of our values as plotted in Figure 2; however, at the minimum near 840  $\text{cm}^{-1}$  the earlier results are lower than ours by 1.5%. At frequencies below 400  $\text{cm}^{-1}$  we have continued satisfactory agreement with the Russians, who based their values in this region on published results of others including *Draeger et al.* [1966] which are shown by the points in Figures 1 and 3 for  $\nu < 200 \text{ cm}^{-1}$ . In the region  $\nu < 400 \text{ cm}^{-1}$  the French results fall below our values and are apparently in serious error; they are based on prism spectrograph results, which we find are subject to stray radiation problems in the low-frequency region.

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